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## ABSTRACT

In this paper, the effect of the Elasto-Hydrodynamic Lubrication (EHL) on the performance of thrust slide-bearings in scroll compressors is considered. The EHL effect accounts for the superior lubrication characteristics of these bearings. The orbiting thrust plate undergoes elastic deformation due to axial loading, resulting in the formation of a uniform fluid wedge between the orbiting and fixed thrust plates. This fluid wedge is a region with very high induced oil film pressure which, in turn, accounts for the remarkably good lubrication characteristics of the thrust slide-bearing. Furthermore, the high oil film pressure further induces a local elastic deformation of the fixed thrust plate to form an EHL pocket in the thrust plate. The greater the deformation, the more the oil film pressure between the sliding surfaces increases, due to the envelopment of the oil in this local region of deformation. The formation of the EHL pocket was verified using FEM analysis and lubrication tests on the elastic deformation of the thrust plate.

**Key Words** : Lubrication, Tribology, EHL, Soft EHL, Thrust Slide-Bearing, Scroll Compressor, Refrigerant

## 1. INTRODUCTION

The scroll compressor is a low-vibration and low-noise high-performance compressor that compresses refrigerant by the use of crescent-shaped multiple compression chambers formed by two pairs of spiral walls of the orbiting and fixed scrolls. Scroll compressors are extensively used in heat pump and air-conditioning applications. Figure 1 shows a cross-sectional view of the compression mechanism of a scroll compressor. Figure 1(a) shows an overall view and Figure 1(b) an enlarged view of the compression mechanism. The lubrication mechanism of the thrust slide bearing of scroll compressor has not yet been definitively elucidated. Recently, however, detailed characterization of the lubrication mechanism has been provided by both experimental and theoretical studies (see Ishii, *et al.* (2008) and Oku, *et al.* (2008)). It was suggested that the elasto-hydrodynamic lubrication (EHL) effect and the envelopment of lubricating oil into EHL pockets on the faces of the orbiting and fixed thrust plates in the thrust slide bearing of a scroll compressor leads to its superior lubrication performance.

In general, elasto-hydrodynamic lubrication (EHL) is understood as a “phenomenon that brings about satisfactory lubricating conditions as the local contact portion makes an elastic deformation on the gear tooth

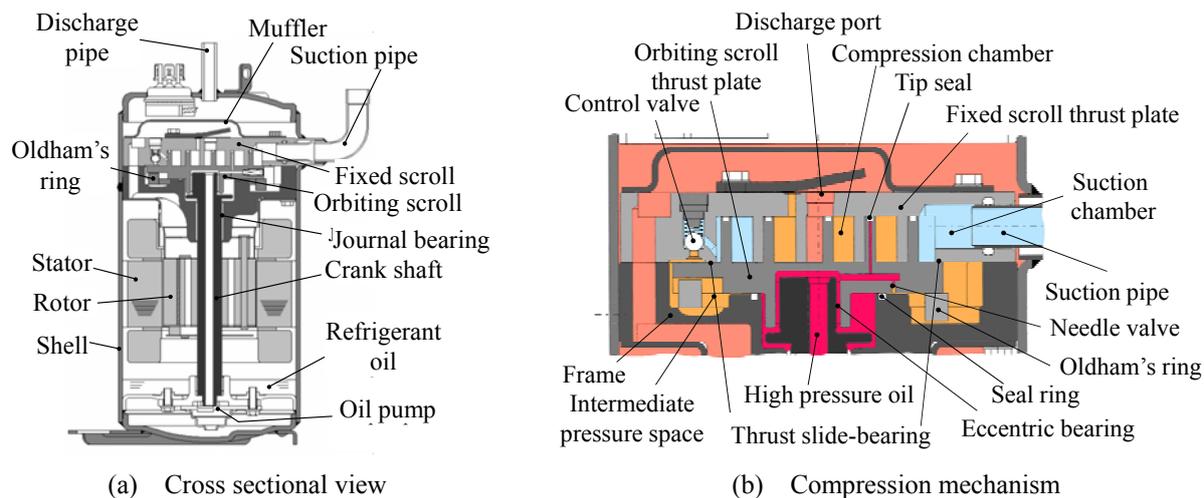


Figure 1 : High pressure type scroll compressor.

contact surface, etc. and in that portion, oil is confined and acquires ultrahigh pressure.” In the thrust slide bearing of scroll compressors, such isolated local contact points do not occur, but the thrust plate undergoes macroscopic elastic deformation due to oil film pressure. The resulting elastic deformation is postulated to give rise to the EHL effect in thrust slide bearings, thereby increasing lubrication performance.

In the present study, an analysis of the fluid lubrication of thrust slide bearing as a mixed lubrication process is presented using the average Reynolds equation of Patir & Cheng (1978 ; 1979) along with the solid contact theory from Greenwood & Williamson. The analysis will specifically addresses the EHL effect in thrust slide bearings, termed **soft EHL**, since it is more dispersed than the EHL previously identified in the lubrication of gear teeth. Further, the analysis does not consider the compressibility of the lubricant nor any changes in its viscosity. The soft EHL presented herein assumes constant lubricant density and viscosity. Subsequently, the elastic deformation of the thrust plate by high oil-film pressure is considered and is shown to give rise to the envelopment of the lubricating oil. Lastly, the formation of pocket-shaped elastic deformations is verified both through experiments on and by FEM analysis of a small-size scroll compressor thrust bearing by Tsuji (2014).

## 2. OUTLINE OF LUBRICATION THEORY ANALYSIS OF THRUST SLIDE BEARING AND SOFT EHL ANALYSIS

The lubrication of thrust slide bearing was theoretically analyzed by Oku, *et al.* (2008). In this section, an outline of their theoretical analysis is presented. The conditions under which the elastic deformation of the thrust plate forms a wedge angle relative around its outer circumference are represented by a theoretical of the rigid body thrust plate, as shown in upper portion of Figure 2. In order to enable the measurement of the friction force working on the orbiting scroll thrust plate shown in Figure 1, the orbiting scroll thrust plate in the model is immobilized, positioned, and supported in the horizontal plane by a pivot bearing positioned on its upper surface. In this manner, the degree of freedom of the thrust plate rotating around the pivot bearing is maintained. Furthermore, since the pivot shaft is allowed to move in the thrust direction, the thrust plate has a degree of freedom in the vertical direction. The orbiting scroll thrust plate in the actual compressor, shown in Figure 1, is loaded by a pressure difference and it deforms, producing a

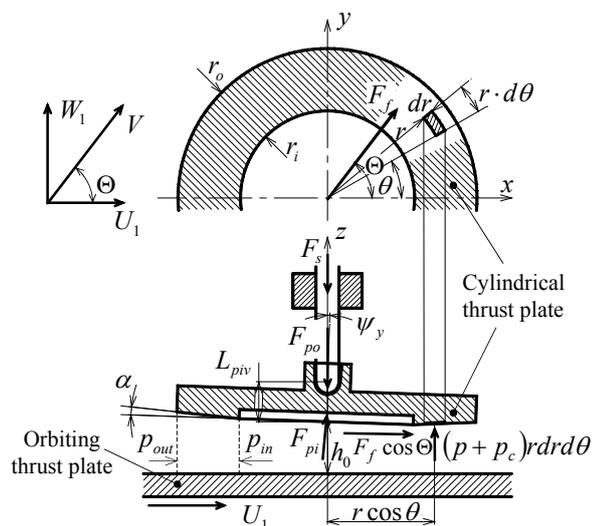


Figure 2 : Mathematical model of rigid thrust slide-bearing with a forced wedge at its periphery, for theoretical analysis of fluid lubrication.

wedge angle around its outer periphery. In the thrust plate model (in Figure 2), the bottom surface of the now fixed equivalent of the orbiting thrust plate is inclined at an angle of  $\alpha$  to provide a geometry relative similar to the elastically deformed orbiting thrust plate. The fixed scroll thrust plate, shown on the top in Figure 1, is located on the lower side of the model and is powered by a motor to undergo an orbiting motion. In this manner, the same relative motion is obtained (for further detail, see Ishii *et al.* (2008) and Oku, *et al.* (2008)). In the model (Figure 2), a large gas force  $F_{po}$  acts downward on the top surface of the cylindrical thrust plate (now fixed) and a small gas force  $F_{pi}$  acts upward on its bottom surface. Due to this difference in pressure loading, the cylindrical thrust plate is pressed against the now orbiting scroll thrust plate. To enable experiments with no gas force difference, a downward acting spring thrust load  $F_s$  is applied to the fixed scroll thrust plate to press it downward. The pivot bearing is located at the center of the fixed scroll thrust plate. The two radii,  $r_i$  and  $r_o$  denote the inner and outer radial, respectively, of the annular region of the thrust slide bearing where the fixed cylindrical thrust plate is pressed against the orbiting scroll thrust plate. The fixed scroll thrust plate can rotate around both the  $x$ -axis and the  $y$ -axis. The symbols  $\psi_x$  and  $\psi_y$  denote the rotation angles around  $x$ -axis and  $y$ -axis of the fixed scroll thrust plate, respectively. An oil film is confined between the fixed scroll thrust plate and the orbiting scroll thrust plate (now on the lower side). The two pressures,  $p_{out}$  and  $p_{in}$ , denote the boundary pressure on the outer and inner circumferences. The boundary velocity of the oil film is given by the velocity  $V$ , which corresponds to the thrust plate velocity as it undergoes orbital motion. The angle  $\Theta$  denotes the direction of  $V$  from the  $x$ -axis. The total frictional force  $F_f$  works in that direction. The symbols  $U_l$  and  $W_l$  denote the  $x$  and  $y$ -components of  $V$ .

## 2.1 Oil Film Pressure and Oil-Film Force

Expressing the oil film thickness  $h$  in polar coordinates  $(r, \theta)$  with its origin at the bearing center  $O$ , we have

$$h(r, \theta) = h_0 + (r - r_i) \tan \alpha - r \cos \theta \cdot \psi_y + r \sin \theta \cdot \psi_x + \delta(r, \theta; t) \quad (r_i < r < r_o) \quad (1)$$

where  $h_0$  denotes the average clearance height and  $\delta$  the elastic deformation rate. Assuming that the oil can be treated as an isothermal incompressible fluid, the oil-film pressure  $p$  that acts on the rough bearing surface can be computed using Equation (2) below, with the average Reynolds expression derived by Patir & Cheng (1978 ; 1979) written in polar coordinates.

$$\begin{aligned} \frac{1}{R} \frac{\partial}{\partial R} \left( \varphi R H^3 \frac{\partial P}{\partial R} \right) + \frac{1}{R^2} \frac{\partial}{\partial \theta} \left( \varphi H^3 \frac{\partial P}{\partial \theta} \right) = \lambda \cdot \frac{1}{R} \left\{ \frac{\partial}{\partial R} (\bar{H}_T \cos(\theta - \Theta)) - \frac{\partial}{\partial \theta} (\bar{H}_T \sin(\theta - \Theta)) \right\} \\ + \lambda \sigma \frac{1}{R} \left\{ \frac{\partial}{\partial R} (\varphi_s R \cos(\theta - \Theta)) - \frac{\partial}{\partial \theta} (\varphi_s \sin(\theta - \Theta)) \right\} + \sigma_s \frac{\partial \bar{H}_T}{\partial \tau_t} \end{aligned} \quad (2)$$

where the non-dimensional variables are defined by the following equations have been introduced.

$$R \equiv r / r_o, \quad P \equiv p / p_a, \quad H \equiv h / h_{ref}, \quad \bar{H}_T \equiv \bar{h}_T / h_{ref}, \quad \tau_t \equiv \omega t \quad (3)$$

where  $p_a$  denotes atmospheric pressure and  $h_{ref}$  an arbitrary standard oil film thickness in the thrust slide bearing, assumed to be 10  $\mu\text{m}$ , as a typical value.  $\bar{h}_T$  represents the average oil film thickness which can be obtained by integrating the local film thickness  $h_T$ .  $\omega$  denotes the orbiting velocity of the orbiting scroll thrust plate (lower side) and  $t$  is time.  $\sigma$  denotes the standard deviation of surface roughness of the composite top and bottom surfaces ( $= \sqrt{\sigma_1^2 + \sigma_2^2}$ ).  $\sigma_1$  and  $\sigma_2$  are standard deviations of the respective asperity heights of the orbiting scroll thrust plate and the fixed scroll thrust plate.  $\lambda$  in Equation (2) denotes the number of bearings and  $\sigma_s$  the squeeze number, which are defined as follows:

$$\lambda \equiv \frac{6\mu^* V}{r_o p_a} \left( \frac{r_o}{h_{ref}} \right)^2, \quad \sigma_s \equiv \frac{12\mu^* \omega}{p_a} \left( \frac{r_o}{h_{ref}} \right)^2 \quad (4)$$

where,  $\mu^*$  denotes the viscosity of lubricating oil. The symbols  $\varphi$  and  $\varphi_s$  denote the pressure-flow coefficient and shear-flow coefficient, respectively. In this model, the thrust plate is assumed to have isotropic roughness.

The oil-film pressure generated on the bearing surface can be found by solving the average Reynolds equation under the boundary conditions. Integrating the oil-film pressure over the whole bearing surface gives the oil-film force  $F_{oil}$  as follows:

$$F_{OIL} = \iint p(r, \theta) r d\theta dr \quad (5)$$

The viscous resistance force  $F_{VS}$  that works on the bearing surface can be calculated by the following equation with the surface roughness taken into account:

$$F_{VS} = \iint \frac{\mu^* V}{h} [(\phi_f + \phi_{fs}) - 2V_{r2}\phi_{fs}] r d\theta dr \quad (6)$$

where  $\phi_f$  and  $\phi_{fs}$  denote shear stress coefficients and  $V_{r2}$  denotes the fractional deviation of roughness.

## 2.2 Solid Contact Force and Solid Frictional Force

According to the solid contact theory of Greenwood & Williamson (1966), the local solid contact ratio  $\alpha^*(r, \theta)$  is given by the following equation:

$$\alpha^*(r, \theta) \left( \equiv \frac{dA}{rd\theta dr} \right) = \pi\eta\beta\sigma \int_{H_r}^{\infty} (s - H_r) \phi^*(s) ds \quad (7)$$

$\phi^*(s)$  represents a distribution function of asperity summit height normalized by the standard deviation of asperity summit height, whereas  $\beta$  denotes the asperity tip curvature radius.  $H_r$  is defined by the ratio of apparent clearance ( $h - 1.1\sigma$ ) to the standard deviation  $\sigma^* (= 0.7\sigma)$  of asperity summit height in the Greenwood-Williamson model and is given by the following equation:

$$H_r = \frac{h - 1.1\sigma}{\sigma^*} \quad (8)$$

If Equation (7) is used, the solid contact force  $F_{sc}$  and the solid frictional force  $F_{ss}$  can be found by the following equation:

$$F_{sc} \left( \equiv \int p_c \cdot dA \right) = \iint p_c \alpha^*(r, \theta) \cdot r d\theta dr, \quad F_{ss} \left( \equiv \int \tau \cdot dA \right) = \iint \tau \alpha^*(r, \theta) \cdot r d\theta dr \quad (9)$$

where  $p_c$  denotes the plastic flow pressure and  $\tau$  the shearing strength.

## 2.3 Bearing Position

The force and moment equilibrium equations are as follows:

$$\left. \begin{aligned} -m\ddot{h}_0 + F_{pi} + F_{OIL} + F_{sc} - F_s - F_{po} &= 0 \\ -I_x \ddot{\psi}_x + \iint p(r, \theta) \cdot r \sin \theta \cdot r dr d\theta + \iint \alpha^*(r, \theta) \cdot p_c \cdot r \sin \theta \cdot r dr d\theta - L_{piv} \cdot F_f \sin \Theta &= 0 \\ -I_y \ddot{\psi}_y - \iint p(r, \theta) \cdot r \cos \theta \cdot r dr d\theta - \iint \alpha^*(r, \theta) \cdot p_c \cdot r \cos \theta \cdot r dr d\theta + L_{piv} \cdot F_f \cos \Theta &= 0 \end{aligned} \right\} \quad (10)$$

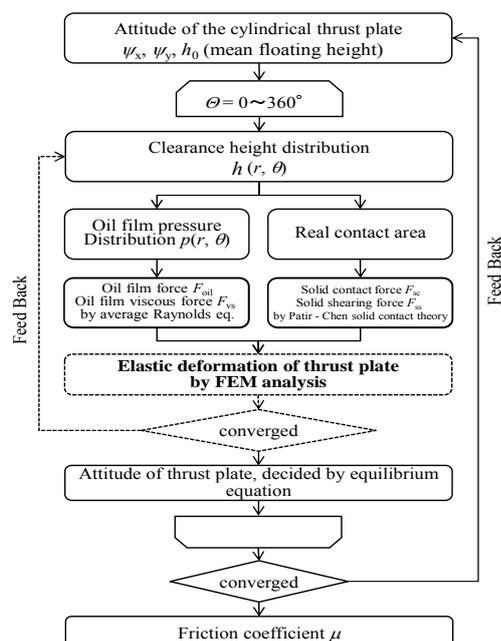
and include: the oil film shear force  $F_{vs}$ ; the solid shear force  $F_{ss}$  applied to the rigid body thrust plate; the solid contact force  $F_{sc}$ ; the oil film force  $F_{OIL}$ ; the spring thrust force  $F_s$ ; the vertical equilibrium between forces  $F_{po}$  and  $F_{pi}$  due to gas pressures inside and outside the bearing; and the moments around the pivot bearing.

The bearing model formulated here assumes complete axisymmetry, and therefore the moments around the z-axis become zero. Consequently, it is not necessary to consider the motion equation around the z-axis. From the equilibrium equations above (Equation 10), the bearing position including the bearing clearance height  $h_0$ , inclination angles ( $\psi_x, \psi_y$ ) can be determined.

## 2.4 Procedure for Computing Numerical Values

Figure 3 shows a flow chart of the theoretical analysis of thrust slide bearing lubrication. In the figure, the dashed line shows the computation made including soft EHL. Initially, the bearing position, including bearing clearance height and gradient  $\psi_x, \psi_y$ , is assumed. Based on the assumption, the average Reynolds equation of Patir & Cheng (1978 ; 1979), Equation (2), is solved by the finite difference method (24 grid lines in the radial direction and 120 grid lines in the circumferential direction), and find the oil film force  $F_{OIL}$  from Equation (5), and the oil film viscous resistance force  $F_{vs}$  from Equation (6) are calculated. At the same time, from Equation (9), the solid contact force  $F_{sc}$  and solid frictional force  $F_{ss}$ , are found. In the event that elastic deformation is taken into account, FEM-analysis of the elastic deformation of the fixed scroll thrust plate by the oil film

pressure is undertaken. Using the FEM deformation, the oil film height distribution and the oil film pressure are computed iteratively until they converge. Finally, the bearing position is then iteratively calculated until the bearing position satisfies Equation (10) and converges.



**Figure 3 :** Computation procedure for lubrication and attitude of thrust slide-bearing, without and with EHL pocket (dashed lines to be added for the case with EHL pocket)

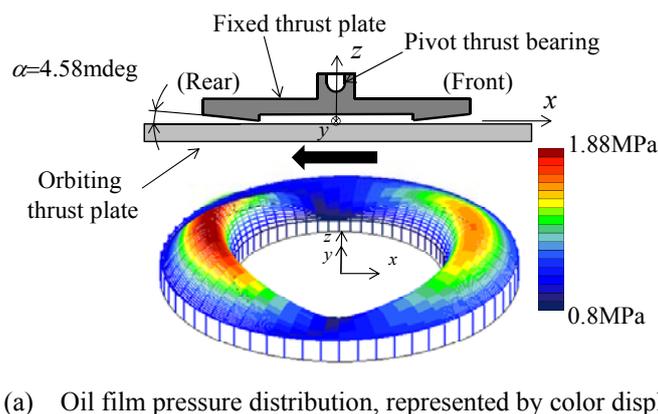
**Table 1 :** Major specifications for computations for rigid thrust slide-bearing

Surface roughness $R_a$	Orbiting thrust plate [ $\mu\text{m}$ ]	3.0
	Cylindrical thrust plate [ $\mu\text{m}$ ]	0.056(in) ~ 0.27(out)
Standard deviation of surface roughness	Orbiting thrust plate $\sigma_1$ [ $\mu\text{m}$ ]	1.458
	Cylindrical thrust plate $\sigma_2$ [ $\mu\text{m}$ ]	0.188(in) ~ 1.15(out)
Bearing dimension	Outer radius $r_o$ [mm]	65.0
	Inner radius $r_i$ [mm]	37.85
	Friction surface area $S_f$ [ $\text{mm}^2$ ]	8772.5
Wedge angle $\alpha$ [mdeg]		2.92
Oil film thickness $h_o$ [mm]		~ 4.64 ~
Thrust plate thickness $T_{th}$ [mm]		10.4
Pivot height $L_{piv}$ [mm]		15.0
Cylindrical thrust plate mass $m$ [kg]		0.340
Moment of inertia $I_x, I_y$ [ $\text{kg}\cdot\text{m}^2$ ]		$3.55 \times 10^{-4}$
Plastic flow pressure $p_c$ [MPa]		1600
Shearing strength $\tau$ [MPa]		240
Surface density of asperities $\eta$ [ $\text{mm}^{-2}$ ]		400
Asperity summits radius $\beta$ [ $\mu\text{m}$ ]		8.0
Oil viscosity $\mu^*$ [Pa·s]		0.00702
Temperature at the surface of thrust plate [ $^{\circ}\text{C}$ ]		70.0
Inner pressure $p_{in}$ [Mpa]		1.0 ~ 0.8 ~ 0.5
Outer pressure $p_{out}$ [MPa]		1.1
Pressure difference $\Delta p$ [MPa]		0.1 ~ 0.3 ~ 0.6
Nominal gas thrust force $F_p$ [kN]		~ 11.0 ~
Spring thrust force $F_s$ [N]		600
Resultant thrust force $F_T$ [kN]		~ 11.6 ~
Orbiting speed $N$ [rpm]		3600
Orbiting radius $r_{orb}$ [mm]		3.0
Sliding velocity $V$ [m/s]		1.13
Number of grids	Radial	24
	Circumferential	120
Young modulus $E$ [GPa]		73.5

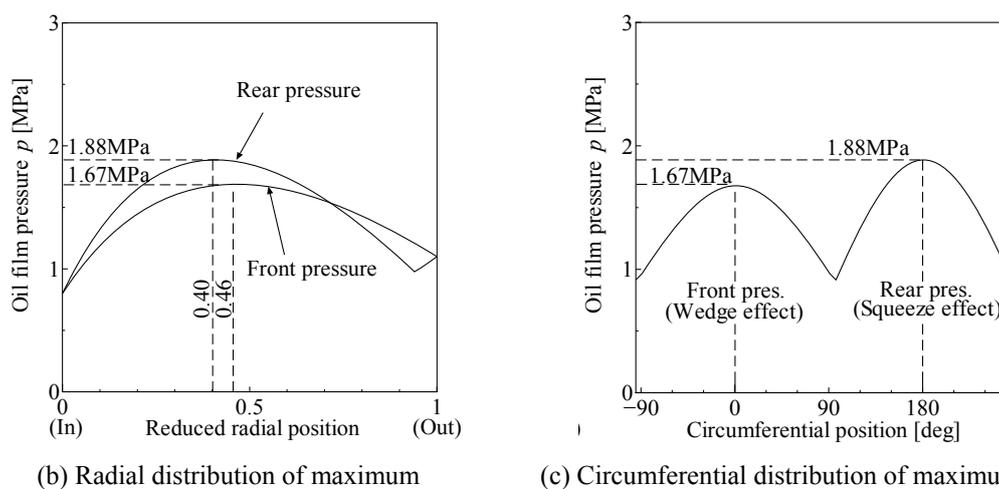
### 3. ENVELOPMENT OF LUBRICATING OIL BY ELASTIC DEFORMATION OF THRUST PLATE

Numerical computation of the oil-film pressure distribution of the annular solid body model neglecting the EHL effect was undertaken by omitting the dashed line portion of the flow chart shown in Figure 3. Figure 4 shows the resulting oil pressure distribution for the model specifications given in Table 1. These specifications correspond a small scroll compressor with the motor shaft input of 0.1 kW. The annular rigid body model had the same area as that of the thrust bearing shown in the top of Figure 4(a). The inner and the outer radii were  $r_i = 37.85$  mm and  $r_o = 65$  mm, respectively. A pressure difference of  $\Delta p = 0.3$  MPa, net downward, was applied across the bearing (between inside and the outside of bearing). In addition, an axial load of  $F_s = 600$  N was applied to the pivot bearing section. The sum of thrust loads  $F_p$  was 11.0 kN. A mean wedge angle of  $\alpha = 2.92$  mdeg for one full rotation was generated on the outer circumference section of the thrust plate by this thrust load based on our FEM analysis of the thrust slide bearing. Subsequently, this wedge angle was used for the outer circumference of the annular rigid body mode. The viscosity coefficient  $\mu^*$  was assumed to be 0.0081 Pa·S with an average temperature of 67.8 $^{\circ}\text{C}$  and average pressure of 0.8 MPa.

With the bottom thrust plate orbiting at a radius of  $r_{orb} = 3.0$  mm and at a rotational speed of  $N = 3600$  rpm, the computed oil-film pressure shown by the contour plot in Figure 4(a) was generated by the wedge effect and the squeeze effect, on the friction surface of thrust slide bearing. This oil-film pressure corresponds to the orbiting scroll thrust plate (bottom plate) moving to the left, that is, in the negative x-direction. Pressure increases from the blue color to the red color. The oil film thickness was 5.3  $\mu\text{m}$  on average and 4.6  $\mu\text{m}$  at minimum. The oil film pressure reaches a maximum along on the x-axis. Figure 4(b) shows the radial distribution of the oil film pressure in terms of the dimensionless radial position. In this plot, the inner circumferential position is zero and the outer circumferential position is 1.0. The oil film pressure distribution is convex upwards with a maximum value of 1.67 MPa at the front section (right side) where the radial position was 0.46 and 188 MPa on the rear section (left side) where the radial position was 0.40. Figure 4(c) shows the circumferential distribution of the maximum oil film pressure in which the angular position is measured



(a) Oil film pressure distribution, represented by color display.



(b) Radial distribution of maximum

(c) Circumferential distribution of maximum pressure

**Figure 4 :** Oil film pressure distribution for  $\alpha = 2.92$  mdeg due to wedge and squeeze effects, without EHL effect.

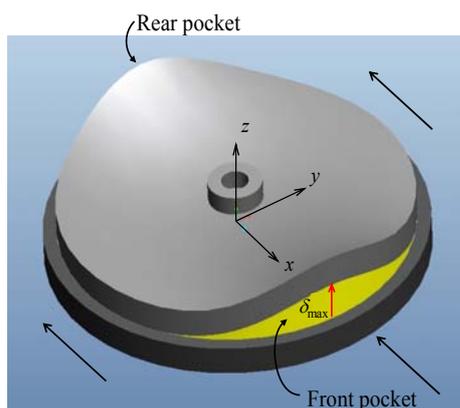
counterclockwise from the x-axis. The oil film pressure on both the front and rear sections reaches a maximum at 0 degrees and +180 degrees (on the x-axis), respectively. The maximum pressure was greater on the rear section. The pressure at the front section, although somewhat smaller occurs at a greater radial distance where its contribution to the moment is greater. The pressure on the front section was primarily generated by the wedge effect, while that on the rear section primarily generated by the squeeze effect.

Since this large oil-film pressure was upwards convex along both the radial and circumferential directions, it was naturally assumed that the thrust plate would also be elastically deformed into an upwards convex shape with this large oil film pressure. Figure 5 illustrates the elastic deformation. Since the thrust plate was loaded downward by the large gas force, the upwards convex elastic deformation tapered off at smaller radii. As expected, the elastic deformation also tapered off in the circumferential direction. This spatial deformation of the frictional surface is referred to as a pocket-shaped region in this paper. As this pocket-shaped region formed, lubricating oil particles were raked into this space as the lower thrust plate moved in the negative x-axis direction and the oil envelopment by the pocket was generated. In this manner, an even greater oil film pressure was generated. As a result, it was expected that an even better lubricating performance would be achieved at the thrust slide bearing section. This phenomenon is called “the EHL effect by lubricating oil envelopment” in this paper.

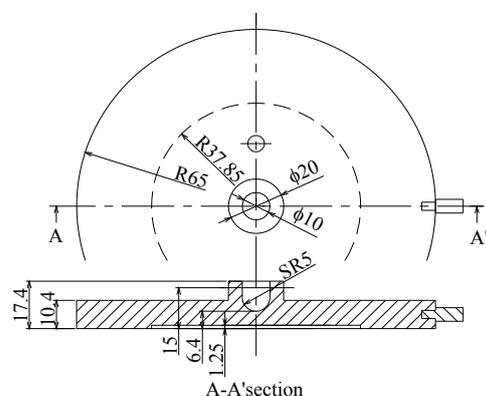
The elastic deformation of thrust plate will be verified both computationally and experimentally in the following section.

#### 4. FEM ANALYSIS AND EXPERIMENTAL VERIFICATION OF EHL POCKET FORMATION

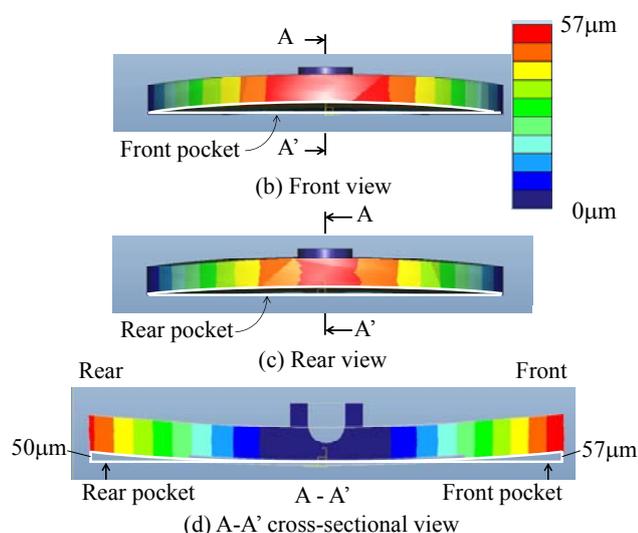
The pocket-shaped region generated on the frictional surface by elastic deformation of the thrust plate is called the EHL pocket in this paper. The elastic deformation of the thrust plate can be validated by FEM analysis. The analysis software, PTC Pro/ENGINEER Wildfire 5.0 (on a Dell DIMENSION 8300; CPU: Core2duo, 6 shows the annular thrust plate along with the computed deformation. Figure 6(a) shows the



**Figure 5 :** Additional elastic deformation of thrust plate due to high oil-film pressure, making EHL pocket.



(a) Cross-sectional view of thrust plate



**Figure 6 :** Thrust plate and its deformation, simulated by FEM analysis

located at the center of thrust plate, the moment was free and movements in the right/left and up/down directions were constrained. Furthermore, 1.1 MPa of pressure was applied to the outside of the bearing, and 0.8 MPa of pressure to the inside. Thus, pressure difference was 0.3 MPa; the thrust gas force due to the pressure difference was 11.0 kN. The oil film pressure shown in Figure 4 was applied to the frictional surface (lower side of the thrust plate) under these conditions.

Figure 6(b) and (c) show the elastic deformation on the front and rear faces of the thrust plate. The deformation increases as the color changes from blue to red. In both views, an upwards-convex deformation was found, with no deformation along the two sides. Consequently, the frictional surface became tapered in the circumferential direction. Figure 6(d) shows section A-A'. Both front and rear surfaces (appearing as the right and left side in Figure 6(d)) were lifted and deformation was maximal on the outer circumferential portion of the clearance. Consequently, the frictional surface became tapered toward the inner circumference, too. The FEM analysis identified the pocket-shaped regions formed at both the front and the rear of the thrust bearing surface (lower side of the thrust plate) by the oil film pressure, as was postulated in the preceding section.

The biggest displacement of the outer circumferential portion of the thrust plate was 57  $\mu\text{m}$  at the front section and 50  $\mu\text{m}$  at the rear section. The thrust plate was substantially deformed by the oil film pressure, but the analysis assumed the oil film pressure remained constant even as the displacement increased. In reality, as the pocket size increases, the oil film pressure decreases. Consequently, the elastic deformation would be correspondingly smaller as the oil film pressure and elastic deformation converge to an appropriate condition where the oil film pressure and elastic deformation would be balanced.

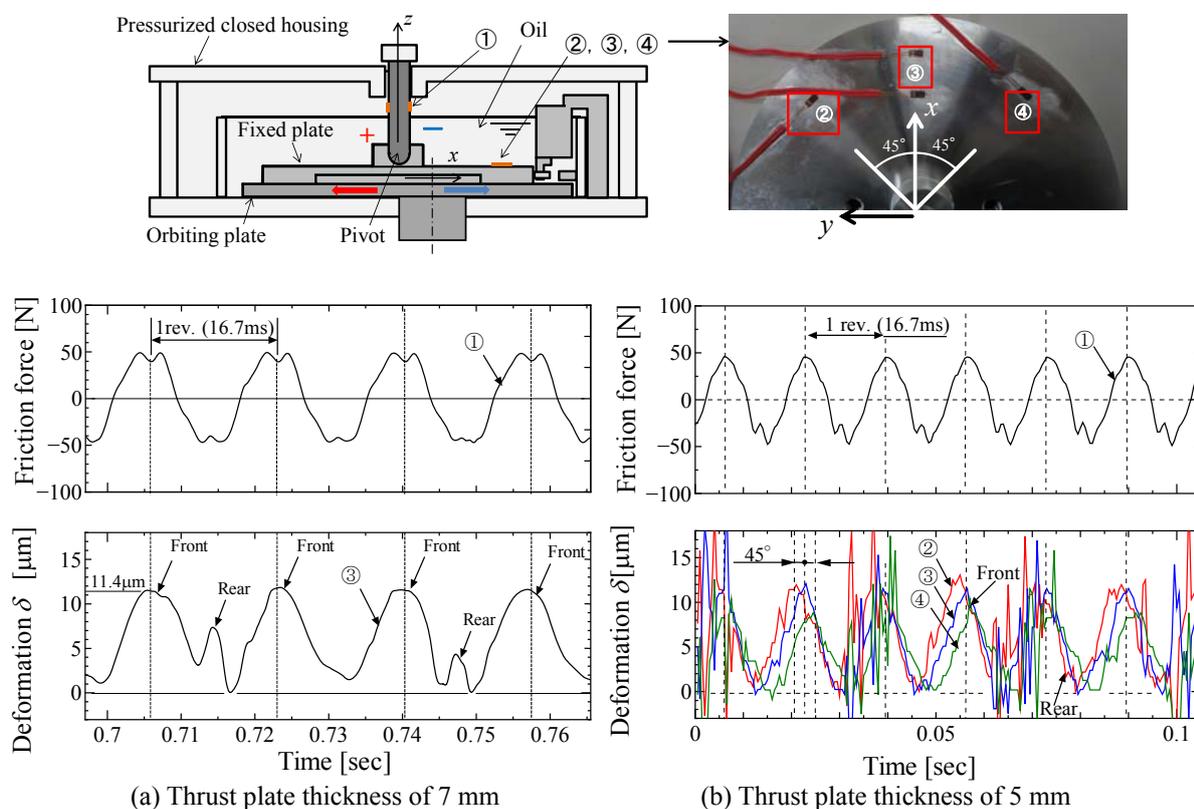
In order to validate further the pocket formation identified in the FEM analysis, friction experiments were conducted in an R-22 atmosphere. The upper portion of Figure 7 shows the set-up for the friction experiments.

The pivot bearing section of the fixed thrust plate (inner radius: 37.85 mm; outer radius: 65 mm) was held by a shaft; however, the moment is not constrained. A swing stopper is attached to the thrust plate, shown in the right-most portion of Figure 7. The lower-side thrust plate was allowed to undergo orbiting motion. The entire thrust slide bearing was immersed in an oil tank and housed inside a pressure vessel in an R-22 atmosphere pressurized to 1.1 MPa. The internal portion of the thrust slide bearing was depressurized to 0.5 MPa. Consequently, the external/internal pressure of the bearing was  $\Delta p = 0.6$  MPa. The lower-side thrust plate underwent orbital motion at  $N = 3600$  rpm.

In order to detect the drag direction of the friction force on the thrust plate, strain gauges [1] were mounted to both sides (on the x-axis) of the shaft. The waveforms in the upper row of Figure 7(a) and (b) are the outputs from the strain gauges. The elapsed time is along the abscissa. The strain was converted into the friction force working on the friction surface using calibration data and is plotted as the ordinate. At the instant when the friction force reaches a positive peak, the lower orbiting scroll thrust plate moves in the negative x-direction. The maximum friction force was about 50 N.

In order to detect elastic deformation due to changes in the oil film pressure acting on the thrust plate, a strain gauge [3] was mounted on the upper right surface of the fixed scroll thrust plate along the x-axis. In addition, strain gauges [2] and [4] were similarly mounted, but displaced  $\pm 45$  degrees relative to the x-axis. The photograph of Figure 7 shows how they were mounted. The relevant radial positions were the center of thrust bearing section. The active gauge direction was the circumferential direction. For an assumed upwards-convex elastic deformation in the circumferential direction, these strain gauges would detect a tensile strain at their maximum sensitivity. To more easily identify pocket formation, the thrust plate thickness was reduced sequentially to 7 mm and then 5 mm, and the experiments were conducted. Since the oil-film pressure distribution shape was not changed, even with reduced plate thickness, the deformation shape appeared in a manner consistent with the FEM analysis results. The lower rows of Figure 7(a) and (b) show the measurements of elastic deformation. The strain was converted to the pocket opening height  $\delta$ , as was shown in Figure 5, using by calibration data, and  $\delta$  is plotted as the ordinate.

The typical height of elastic deformation  $\delta$  indicated a peak precisely synchronized with the friction force. This variation of  $\delta$  did not arise from the frictional force. Even if the 50 N-friction-force worked on the surface below the thrust plate, the effect could be scarcely sensed as the strain orthogonal to the frictional force in the



**Figure 7 :** Measured dynamic deformation of the thrust plate, operated at 3600 rpm and  $\Delta p = 0.6$  MPa.

circumferential direction on the top surface of thrust plate. This observation was confirmed using the FEM analysis, which predicted a circumferential strain due to the 50 N frictional force of about  $0.8 \mu\epsilon$  for a thrust plate thickness  $t = 5.0$  mm. The  $0.8 \mu\epsilon$  was only about 1.3% of the strain rate generated by the oil film force. Other forces exerted on the thrust plate included the gas pressure and axial load by spring force, but these were held constant and did not cause variations in strain. At the same time, even if the axial load was varied, it would not result in tensile strain in the circumferential direction on the top surface of thrust plate, rather a microscopic compression strain would be generated. The compression strain would amount to only about 0.1% of the strain rate generated by the oil film force even if the axial load was assumed to vary by 0.1%. Consequently, the axial load was unconnected to variations of  $\delta$ .

Among other factors that might be assumed to give rise to the elastic deformation  $\delta$ , one might consider the force due to oil film pressure, shown in Figure 4, generated on the frictional surface. When  $\delta$  reaches its peak, the lower orbiting scroll thrust plate moves to the left (negative x axis), and therefore, the front portion of the orbiting plate is directly beneath the location of strain gauge [3]. Since a wedge oil film pressure of up to 3 MPa was generated at the front portion of the plate, it would be appropriate to suggest that the thrust plate would undergo elastic deformation in the upward direction. The maximum typical height  $\delta$  was about  $11 \mu\text{m}$  when the thrust plate was 7 mm thick (Figure 7(a)). Obviously, as the thrust plate orbited, the wedge oil film pressure rotated over friction surface, and therefore, the oil film pressure right below the strain gauge gradually decreased and the measured height  $\delta$  decreased accordingly. This result suggests that the thrust plate underwent upwards-convex deformation in the circumferential direction as shown in Figure 5 and Figure 6 when the deformed shape at certain times in the rotation cycle.

Near the negative peak in the friction force, that is, near where the orbiting scroll thrust plate was moving in the right (x-axis), another increase in the deformation height  $\delta$  was again observed. At this time, the rear section was located directly beneath the strain gauge and a large squeeze oil film pressure (Figure 4) was generated. This variation, assume to be generated by the large squeeze oil film pressure failed to appear in the following cycle, but then re-appeared in the subsequent cycle. The variation tended to appear once every two rotations and its amplitude was small, as well. Since the squeeze pressure was expected to decrease due to elastic deformation, it was assumed that the squeeze pressure might have given rise to this kind of unstable variation. Figure 7(b) shows the experimental results when the thrust plate thickness was further reduced to 5 mm. At the rear section, severe chattering was produced. Although we cannot make any definitive statement at the present stage, one might assume that when the plate thickness was reduced, the plate was more easily bent, and the elastic deformation and the squeeze pressure were both strong and in opposite phase, resulting in severe chattering. This figure shows that the outputs of the strain gauges [2] and [4] at  $\pm 45$  degrees also exhibited similar chattering. The obvious phase difference of elastic deformation output that appeared on the front section was 45 degrees, consistent with the offset angle at which the strain gauges were mounted, indicating that the pocket formed at the front section was synchronized with the orbiting motion.

## 5. CONCLUSION

Studies were made by experiment and theoretical analysis on the possibility of superb lubrication characteristics brought about by an EHL pocket which is formed by a thrust plate that locally undergoes elastic deformation by large oil film pressure generated in the thrust slide bearing section of a scroll compressor. The studies conducted using thrust plates of dimensions appropriate for a small-size scroll compressor with a motor shaft output of 0.1 kW. The results are summarized as follows:

- (1) FEM analysis was undertaken to investigate the elastic deformation of the thrust plate by the oil film pressure; it was confirmed that a pocket-shaped region was formed.
- (2) In an R-22 atmosphere, friction experiments were conducted which confirmed large elastic deformation in the front section of the thrust plate due to wedge oil film pressure; the position rotated in synchronism with the orbiting movement. In addition, the rear section was observed to undergo smaller elastic deformation relative to that of the front section which was generated once every two orbiting rotations and severe chattering occurred when the thrust plate was thin.

Discussion of the oil film pressure increase due to the oil envelopment are considered in our companion paper, Ishii, *et al.* (2014), where further studies on the EHL effect, including the calculation of the coefficient of friction in the thrust slide-bearing with and without EHL oil envelopment is presented.

## NOMENCLATURE

$dA$	: Local real contact area, $m^2$	$V_{r1}, V_{r2}$	: Variance ratio, -
$F_f$	: Resultant frictional force, N	$\alpha$	: Wedge angle, rad
$F_{OLL}$	: Oil film force, N	$\alpha^*$	: Local solid contact ratio, -
$F_p$	: Nominal gas thrust force, N	$\beta$	: Asperity summits radius, m
$F_{pi}, F_{po}$	: Inner gas thrust force, N	$\Delta p$	: Pressure difference, Pa
$F_s$	: Axial spring force, N	$\eta$	: Surface density of asperities, $m^{-2}$
$F_{sc}$	: Solid contact force, N	$\Theta$	: Orbiting angle, rad
$F_{ss}$	: Solid shearing force, N	$\lambda$	: Bearing number, -
$F_{vs}$	: Oil viscous force, N	$\lambda^*$	: Non-dimensional factor for oil flow velocity, -
$h$	: Nominal oil film thickness, m	$\mu_{th}$	: Coefficient of friction, -
$h_0$	: Average clearance height, m	$\mu^*$	: Oil viscosity, Pa · s
$h_{ref}$	: Arbitrary standard thickness, m	$\sigma, \sigma_1, \sigma_2$	: Standard deviations of surface combined roughness, m
$H_r$	: Oil film thickness to surface roughness ratio, -	$\bar{\sigma}$	: Non-dimensional deviations of surface combined roughness, -
$\bar{h}_r$	: Average oil film thickness, m	$\sigma_s$	: Squeeze number, -
$N$	: Orbital speed, rpm	$\tau_i$	: Non-dimensional time, rad
$p$	: Oil film pressure, Pa	$\tau$	: Shearing strength, Pa
$p_a$	: Atmospheric pressure, Pa	$\phi$	: Pressure flow factor, -
$p_c$	: Plastic flow pressure, Pa	$\phi_s$	: Shear flow factor, -
$p_{in}, p_{out}$	: boundary pressure, Pa	$\phi_f, \phi_s$	: Shear stress factor, -
$r_i, r_o$	: Inner bearing radius, m	$\psi_x, \psi_y$	: Rotation angle rotating $x$ - axis, rad
$\bar{U}, \bar{W}$	: Non-dimensional average oil flow velocity, -	$\omega$	: Angular orbiting velocity, rad/s
$\bar{u}, \bar{w}$	: Average oil flow velocity, m/s		
$U_l, V, W_l$	: Boundary velocity, m/s		

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