Measurement of the charge asymmetry in 
$B \rightarrow K^{-*}(892)(+/+)\pi(-/+)$

Measurement of the charge asymmetry in $B \to K^*(892)^{\pm} \pi^{\mp}$

University of Illinois, Urbana-Champaign, Illinois 61801, USA

K. W. Edwards
Carleton University, Ottawa, Ontario, Canada K1S 5B6
and the Institute of Particle Physics, Canada

D. Besson
University of Kansas, Lawrence, Kansas 66045, USA

University of Minnesota, Minneapolis, Minnesota 55455, USA

Z. Metreveli, K.K. Seth, A. Tomaradze, and P. Zweber
Northwestern University, Evanston, Illinois 60208, USA

S. Ahmed, M. S. Alam, J. Ernst, L. Jian, M. Saleem, and F. Wappler
State University of New York at Albany, Albany, New York 12222, USA

Ohio State University, Columbus, Ohio 43210, USA

H. Severini and P. Skubic
University of Oklahoma, Norman, Oklahoma 73019, USA

S.A. Dytman, J.A. Mueller, S. Nam, and V. Savinov
University of Pittsburgh, Pittsburgh, Pennsylvania 15260, USA

Purdue University, West Lafayette, Indiana 47907, USA

D. Cronin-Hennessy, C. S. Park, W. Park, J. B. Thayer, and E. H. Thorndike
University of Rochester, Rochester, New York 14627, USA

T. E. Coan, Y. S. Gao, F. Liu, and R. Stroynowski
Southern Methodist University, Dallas, Texas 75275, USA

Syracuse University, Syracuse, New York 13244, USA

A. H. Mahmood
University of Texas–Pan American, Edinburg, Texas 78539, USA

S. E. Csorna and I. Danko
Vanderbilt University, Nashville, Tennessee 37235, USA

G. Bonvicini, D. Cinabro, M. Dubrovin, and S. McGee
Wayne State University, Detroit, Michigan 48202, USA

A. Bornheim, E. Lipeles, S. P. Pappas, A. Shapiro, W. M. Sun, and A. J. Weinstein
California Institute of Technology, Pasadena, California 91125, USA

R. A. Briere, G. P. Chen, T. Ferguson, G. Tatishvili, H. Vogel, and M. E. Watkins
Carnegie Mellon University, Pittsburgh, Pennsylvania 15213, USA
Cornell University, Ithaca, New York 14853, USA

S. B. Athar, P. Avery, L. Breva-Newell, V. Potlia, H. Stoeck, and J. Yelton
University of Florida, Gainesville, Florida 32611, USA

(CLEO Collaboration)

(Received 22 April 2003; published 10 July 2003)

We report on a search for a CP-violating asymmetry in the charmless hadronic decay \( B \to K^0(892)^\pm \pi^\mp \), using 9.1 fb\(^{-1}\) of integrated luminosity produced at \( \sqrt{s} = 10.58 \) GeV and collected with the CLEO detector. We find \( A_{CP}(B \to K^0(892)^\pm \pi^\mp) = 0.26^{+0.33}_{-0.30} \text{(stat)}^{+0.10}_{-0.08} \text{(syst)} \), giving an allowed interval of \([-0.31, 0.78]\) at the 90% confidence level.

DOI: 10.1103/PhysRevD.68.017101

PACS number(s): 13.20.He

The standard model predicts that CP-violating phenomena are governed solely by the single imaginary parameter of the Cabibbo-Kobayashi-Maskawa matrix \([1]\) of complex quark couplings. The first observations of CP-violation in the neutral B system were recently reported \([2]\), and they have been interpreted widely as induced by \( B^0\overline{B}^0 \) mixing. To date, direct CP violation has been observed only in the neutral kaon system \([3]\). Direct CP violation in a given decay requires contributions from two or more amplitudes which differ in both CP-violating (weak) and CP-conserving (strong) phases. In the B system, these conditions are expected to be met in some charmless hadronic decays, and direct CP violation can occur at sizable levels, depending on the magnitude of the strong phase difference \([4,5]\) or on the presence of new physics \([6]\). Previous analyses, mainly focusing on two-pseudoscalar final states, have not observed direct CP violation in these decays \([7,8]\). In this Brief Report, we present a search for direct CP violation in the vector-pseudoscalar decay \( B \to K^*(892)^\pm \pi^\mp \). We express the difference between the decay rates for \( B^0 \to K^0(892)^\mp \pi^\pm \) and \( B^0 \to K^* \pi^\mp \) in terms of an asymmetry \( A_{CP} \) defined as

\[
A_{CP} = \frac{B(B^0 \to K^0(892)^\mp \pi^\pm) - B(B^0 \to K^*(892)^\mp \pi^\pm)}{B(B^0 \to K^0(892)^\mp \pi^\mp) + B(B^0 \to K^*(892)^\mp \pi^\pm)},
\]

We consider both \( K^*(892)^\pm \) submodes, \( K^0(892)^\pm \to K_S^0 \pi^\pm \) and \( K^*(892)^\pm \to K^+ \pi^0 \), by analyzing the final states \( K_S^0 \pi^\pm h^\mp \) and \( K^+ h^\pm \pi^\mp \), where \( h^\pm \) denotes a charged pion or kaon. We perform a maximum likelihood fit in the \( K^0_S \pi^\pm h^\mp \) and \( K^+ h^\pm \pi^\mp \) Dalitz plots to distinguish \( B \to K^*(892)^\pm \pi^\mp \) from other intermediate resonances or nonresonant three-body decays. The CP-averaged branching fraction for \( B \to K^*(892)^\pm \pi^\mp \) has been measured by the Belle \([9]\) and CLEO \([10]\) Collaborations, and the work described in this Brief Report is an extension of that previous CLEO analysis.

The data sample used in this analysis was produced in symmetric \( e^+e^- \) collisions at the Cornell Electron Storage Ring (CESR) and collected with the CLEO detector in two configurations, known as CLEO II \([11]\) and CLEO II.V \([12]\). It comprises 9.1 fb\(^{-1}\) of integrated luminosity collected on the \( Y(4S) \) resonance, corresponding to \( 9.7 \times 10^6 \) \( B \overline{B} \) pairs, of which \( 6.3 \times 10^6 \) were taken with CLEO II.V. An additional 4.4 fb\(^{-1}\) collected below the \( B \overline{B} \) production threshold is used to study non-\( B \overline{B} \) backgrounds. Of this latter luminosity, 2.8 fb\(^{-1}\) were collected with CLEO II.V. The response of the experimental apparatus is studied with a detailed GEANT-based \([13]\) simulation of the CLEO detector, where the simulated events are processed in a fashion similar to data.

In CLEO II, the momenta of charged particles are measured with a tracking system consisting of a six-layer straw tube chamber, a ten-layer precision drift chamber, and a 51-layer main drift chamber, all operating inside a 1.5 T superconducting solenoid. The main drift chamber also provides a measurement of specific ionization energy loss (\( dE/dx \)), which is used for particle identification. For CLEO II.V, the six-layer straw tube chamber was replaced by a three-layer double-sided silicon vertex detector, and the gas in the main drift chamber was changed from an argon-helium to a helium-propane mixture. Photons are detected with a 7800-crystal CsI electromagnetic calorimeter, which is also inside the solenoid. Proportional chambers placed at various depths within the steel return yoke of the magnet identify muons.

Charged tracks are required to be well measured and to satisfy criteria based on the track fit quality. They must also be consistent with coming from the interaction point in three dimensions. Pions and kaons are identified by consistency with the expected \( dE/dx \), and tracks that are positively identified as electrons or muons are not allowed to form the \( B \) candidate. We form \( \pi^0 \) candidates from pairs of photons with invariant mass within 20 MeV/c\(^2\) \([\text{approximately } 2.5 \text{ standard deviations} (\sigma)]\) of the known \( \pi^0 \) mass. These candidates are then kinematically fitted with their masses con-
strained to the known $\pi^0$ mass. We also require the $\pi^0$ momentum to be greater than 1 GeV/$c$ to reduce combinatoric background from low-momentum $\pi^0$ candidates. $K_S^0$ candidates are selected from pairs of tracks with invariant mass within 10 MeV/$c^2$ (approximately 2.5$\sigma$) of the known $K_S^0$ mass. In addition, $K^0_S$ candidates are required to originate from the beam spot and to have well-measured displaced decay vertices.

We identify $B$ meson candidates by their invariant mass and the total energy of their decay products. We calculate a beam-constrained mass by substituting the beam energy ($E_b$) for the measured $B$ candidate energy: $M = \sqrt{E_b^2 - p_B^2}$, where $p_B$ is the $B$ candidate momentum. Performing this substitution improves the resolution of $M$ by one order of magnitude, to about 3 MeV/$c^2$. We define $\Delta E = E_1 + E_2 + E_3 - E_b$, where $E_1$, $E_2$, and $E_3$ are the energies of the $B$ candidate daughters. For final states with $K^0_S$ and two charged tracks, the $\Delta E$ resolution is about 20 MeV for CLEO II and 15 MeV for CLEO II. A $\pi^0$ in the final state degrades this resolution by approximately a factor of 2. $\Delta E$ is always calculated assuming the $h^+$ is a pion. Therefore, the $\Delta E$ distribution for pions is centered at zero, while that for kaons is shifted by at least $-40$ MeV. We accept $B$ candidates with $M$ between 5.2 and 5.3 GeV/$c^2$ and with $|\Delta E|$ less than 300 MeV for $K^+ h^- \pi^0$ and 200 MeV for $K^0_S \pi^+ h^-$. This region includes the signal region and a high-statistics sideband for background normalization. We reject candidates that are consistent with the exclusive $b \rightarrow c$ transitions $B \rightarrow D \pi$, where $D \rightarrow K \pi$, and $B \rightarrow \phi K^0$, where $\phi \rightarrow \mu^+ \mu^-$ and the muons are misidentified as pions.

The main background in this analysis arises from $e^+ e^- \rightarrow q \bar{q}$, where $q = u, d, s, c$. To suppress this background, we calculate the angle $\theta_{sph}$ between the sphericity axis [14] of the tracks and showers forming the $B$ candidate and that of the remainder of the event. Because of their two-jet structure, continuum $q \bar{q}$ events peak strongly at $|\cos \theta_{sph}| = 1$, while the more isotropic $B \bar{B}$ events are nearly flat in this variable. By requiring $|\cos \theta_{sph}| < 0.8$, we reject 83% of the continuum background while retaining 83% of signal $B$ decays. Additional separation of signal from $q \bar{q}$ background is provided by a Fisher discriminant [15] $F$ formed from 11 variables: the angle between the sphericity axis of the candidate and the beam axis, the ratio of Fox-Wolfram moments $H_2/H_0$ [16], and the scalar sum of the visible momentum in nine $10^\circ$ angular bins around the candidate sphericity axis. In the likelihood fit, we also make use of the angle between the $B$ candidate momentum and the beam axis, $\theta_B$. Angular momentum conservation causes $B$ mesons produced through the $Y(4S)$ to exhibit a sin$^2\theta_B$ dependence, while candidates from continuum are flat in cos $\theta_B$.

In both the $K^0_S \pi^+ h^-$ and $K^+ h^- \pi^0$ topologies, the $h^-$ refers to the faster of the two tracks, which typically has momentum above 1 GeV/$c$. Because $dE/dx$ still provides limited separation of pions and kaons above 1 GeV/$c$, we make use of the $dE/dx$ information in the likelihood fit. In $B \rightarrow K^+ (892)^\pm \pi^\mp$ decays, this higher-momentum track is the one that recoils from the $K^+ (892)^\pm$, more than 99.99% of the time. Its charge uniquely distinguishes $\bar{B}^0 \rightarrow K^+ (892)^\mp \pi^\pm$ from $B^0 \rightarrow K^+ (892)^\mp \pi^\pm$. Thus, the charge asymmetry $A_{+\mp}$, formed using the charge of this higher-momentum track, is essentially the same as $A_{CP}$.

Our loose selection criteria result in samples consisting primarily of background events and containing 11893 candidates for $K_S^0 \pi^+ h^-$ and 28589 for $K^+ h^- \pi^0$. To extract yields and $CP$ asymmetries, we perform an unbinned maximum likelihood fit using the observables $M$, $\Delta E$, $F$, $\cos \theta_B$, the two Dalitz plot variables in each topology, and the $dE/dx$ of the $h^-$ (the faster of the two primary tracks). At high momentum, charged pions and kaons are statistically separated by their $dE/dx$ and by $\Delta E$, each of which provides discrimination at the 2.0$\sigma$ level (1.7$\sigma$ for CLEO II), and we fit for both $\pi$ and $K$ hypotheses simultaneously. Charged pions and kaons with momentum below 1 GeV/$c$ are cleanly identified by $dE/dx$ consistency at the 3$\sigma$ level. The free parameters in the fit are yields ($N$) summed over charge states $N_{h^+} + N_{h^-}$ and charge asymmetries $A_{+\mp} = (N_{h^+} - N_{h^-})/(N_{h^+} + N_{h^-})$.

The probability for a candidate to be consistent with a given component is the product of the probability density function (PDF) values for each of the input variables (neglecting correlations). The likelihood for each candidate is the sum of probabilities over the components in the fit, with relative weights determined by maximizing the total likelihood of the sample, which is given by the following expression:

$$L = \prod_{i=1}^{\text{candidates}} \left\{ \sum_{j=1}^{\text{components}} \left[ f_j (1 \pm A^i_{+\mp}) \prod_{k=1}^{\text{variables}} \mathcal{P}_{ijk} \right] \right\},$$

where the $\pm$ refers to the charge of $h^\mp$ in each candidate. The $\mathcal{P}_{ijk}$ are the per candidate PDF values, and the $f_j$ and $A^i_{+\mp}$ are the free parameters optimized by the fit. The products $f_j (1 \pm A^i_{+\mp})/2$ are constrained to sum to the fraction of candidates in the fit with the appropriate charge of $h^\mp$. Since the PDFs are normalized to unit integral over the fit domain, the $f_j$ can be interpreted as component fractions. The parameters of the $dE/dx$ PDFs are measured from $D \rightarrow K^- \pi^0$ decays in data. For all other variables, the signal and the background $b \rightarrow c$ PDFs are determined from high-statistics Monte Carlo samples, and the continuum PDFs are determined from data collected below the $B \bar{B}$ production threshold. The impact of correlations among the input variables is reduced by determining the PDFs as a function of the event location in the Dalitz plot, for coarse bins in the $M^2(K\pi) - M^2(\pi\pi)$ plane. We use Monte Carlo simulation to estimate the systematic uncertainty associated with neglecting any remaining correlations.

Events from $B \rightarrow K_S^0 \pi^+ h^-$ and $B \rightarrow K^+ h^- \pi^0$, including $B \rightarrow K^* (892)^\pm \pi^\mp$, are modeled in the fit as follows. We consider various $B$ decay channels with intermediate resonances $[K^*(892), K^0_S (1430), \rho (770), \text{and } f_0 (980)]$ as well as nonresonant phase space decay. The Dalitz plot PDFs include our knowledge of the helicity structure in these decays. We neglect interference among these signal processes and assign a systematic uncertainty estimated from Monte Carlo simulation. The decays $B \rightarrow K^- h^+$, where $K^-$ denotes $K^*(892)^-$ or $K^0_S (1430)^-$, are accessible through different
We perform the fit with differing combinations of intermediate resonant and nonresonant states, with up to 12 signal components. The fitted value of $A_{CP}$ does not depend heavily on the number of signal components in the fit, and we include a systematic uncertainty for these variations. We also allow for four background components: pion and kaon hypotheses for $h^\pm$ for continuum background and background from $b\rightarrow c$ decays. We do not fit for charge asymmetries in the background components, but we measure them to be consistent with zero. The $B \rightarrow K^*(892)^{\pm} \pi^{\mp}$ event yields were measured to be $10^{12.6_{-3.0}^{+4.6}}$ for $K^*(892)^{\pm} \rightarrow K_S^{\pm} \pi^\mp$ and $6.1_{-1.9}^{+2.2}$ for $K^*(892)^{\pm} \rightarrow K^0 \pi^\mp$ with a combined statistical significance of 4.6$\sigma$. In the fit, these yields are corrected for efficiency and cross feed from other modes, and the CP asymmetry in $B \rightarrow K^*(892)^{\pm} \pi^{\mp}$ is measured to be $A_{CP} = 0.26^{+0.34+0.10} -0.34 -0.08$, where the uncertainties are statistical and systematic, respectively. The dominant contributions to the latter are statistical uncertainties in the PDFs and variations in the fitting method.

We determine the dependence of the likelihood function on $A_{CP}$ by repeating the fit at several fixed values of $A_{CP}$. By convoluting this function with the systematic uncertainties and integrating the resultant curve in the physical region, we construct a 90% confidence level interval of $0.31 < A_{CP} < 0.78$, where the excluded regions on both sides each contain 5% of the integrated area. Figure 1 shows the likelihood function given by the fit and the effect of including systematic uncertainties.

In summary, we have measured the CP asymmetry in $B \rightarrow K^*(892)^{\pm} \pi^{\mp}$ using a simultaneous maximum likelihood fit to the $B \rightarrow K_S^{\pm} \pi^\mp$ and $B \rightarrow K^\pm h^\mp \pi^0$ topologies. We obtain the value $A_{CP} = 0.26^{+0.34+0.10} -0.34 -0.08$, which is consistent with the theoretical predictions [5] of $-0.19$ to 0.47. We also establish a 90% confidence level interval of $[-0.31, 0.78]$.

We gratefully acknowledge the effort of the CESR staff in providing us with excellent luminosity and running conditions. M.S. thanks the Research Corporation, and A.H.M. thanks the Texas Advanced Research Program. This work was supported by the National Science Foundation, and the U.S. Department of Energy.


